University of Ibn-Khaldoun, Tiaret Faculty of Matter Sciences

1st year master Medical Physics Solution of the exam: Physics of ionised media Year 2023/2024

Ex 01: (5pt)

The following relation can be used:

$$\omega_{\rm p} = \sqrt{\frac{n_{\rm e} \, e^2}{\varepsilon_0 \, m_{\rm e}}} \tag{1.5pt}$$

$$\lambda_{\rm D} = \sqrt{\frac{T_{\rm e}}{m_{\rm e}}} \,\omega_{\rm p}^{-1} = \sqrt{\frac{\varepsilon_0 \,T_{\rm e}}{n_{\rm e} \,e^2}} \tag{1.5pt}$$

For a cell of a typical plasma display:

$$n_e = 1.66 \ 10^{17} \text{ cm}^{-3}$$
 and $T_e = 1.32 \ \text{eV} \ (1.53 \ 10^4 \ \text{K})$ (0.5pt)+ (0.5pt)

For a fluorescent lamp:

$$n_e = 10^{10} \text{ cm}^{-3}$$
 and $T_e = 1 \text{ eV}$ (1.16 10⁴ K) (0.5pt)+ (0.5pt)

Ex 02: (5pt)

Replacing the Debye-Huckel potential into the Poisson equation and calculate Laplace operator in spherical coordinates

$$\nabla^2 \varphi(r) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \varphi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial \varphi}{\partial \theta} \right) + \frac{1}{r^2 \sin \theta^2} \frac{\partial^2 \varphi}{\partial \phi^2}$$
(1) (1pt)

Replacing Debye-Huckel potential $\varphi(r) = \frac{e}{4\pi\varepsilon_0} \frac{exp\left(\frac{r}{\lambda_D}\right)}{r}$ in equation (1)

gives $\nabla^2 \varphi(r) = \frac{\varphi(r)}{r_D^2}$ (3) where $r_D = \lambda_D$ (2pt)

- Equation (3) is Poisson equation $\nabla^2 \varphi(r) = -\frac{\rho}{\varepsilon_0}$ therefore the charge density is

expressed as
$$\rho = -\varepsilon_0 \frac{\varphi(r)}{r_D^2} = -\frac{e}{4\pi r_D^2} \frac{exp\left(-\frac{r}{\lambda_D}\right)}{r}$$
 (2pt)

Ex 03: (4pt)

The drift velocity is given by the relation $\vec{v}_D = \frac{\vec{E} \wedge \vec{B}}{B^2} = = v_D = \frac{E}{B} = 10^6 \text{m/s}$ (2pt) (2pt)

Ex 04: (6pt)

The cold plasma hypothesis means that the electrons only move under the influence of the field $E = E_0 e^{i\omega t}$, so that:

$$\boldsymbol{v} = \boldsymbol{v}_{\mathbf{0}} e^{i\omega t}$$
 (1pt)

Substituting this solution into the equation

$$m_e \frac{\mathrm{d}\boldsymbol{v}}{\mathrm{d}t} = -e\boldsymbol{E} - m_e \nu \boldsymbol{v}$$

gives
$$\boldsymbol{v} = -\frac{e\boldsymbol{E}_0}{m_e} \frac{1}{\nu + i\omega} \mathrm{e}^{i\omega t} \tag{1pt}$$

By definition, the current density can be written $= -e n v = \sigma E$, from which (0.5pt)

$$\sigma = \frac{ne^2}{m_e} \frac{1}{\nu + i\omega} \tag{0.5pt}$$

- To develop equation wave equations for the oscillating electric and magnetic field, we first consider Faraday's law and Ampere's law

$$\nabla \times \mathbf{E} = \frac{i\omega}{c} \mathbf{B}, \qquad (1) \qquad (0.5\text{pt})$$
$$\nabla \times \mathbf{B} = \frac{4\pi}{c} \sigma \mathbf{E} - \frac{i\omega}{c} \mathbf{E}. \qquad (2) \qquad (0.5\text{pt})$$

Eq(2) can be written

$$\nabla \times \mathbf{B} = -\frac{i\omega}{c} \epsilon \mathbf{E}$$
 (3) (0.5pt)

Where $\varepsilon = 1 + i \frac{4\pi\sigma}{\omega}$ define the dielectric function of the plasma Using equation (1) in eq(3) gives,

$$\nabla^2 \mathbf{E} - \nabla (\nabla \cdot \mathbf{E}) + \frac{\omega^2}{c^2} \epsilon \mathbf{E} = \mathbf{0}. \qquad (0.5 \text{pt}) \qquad (4)$$

Since $\nabla \cdot \mathbf{E} = \mathbf{0}$ and assuming a spatial dependence of $\exp(i\mathbf{k} \cdot \mathbf{x})$ then eq (4) and

gives the dispersion relation for electromagnetic waves in plasma

$$\frac{\omega^2}{c^2} \epsilon = k^2 \tag{5} \tag{0.5pt}$$